

Cross section measurements of high- p_T dilepton final-state processes using a global fitting method

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We present a new method for studying high- p_T dilepton events ($e^\pm e^\mp$, $\mu^\pm \mu^\mp$, $e^\pm \mu^\mp$) and simultaneously extracting the production cross sections of $p\bar{p} \rightarrow t\bar{t}$, $p\bar{p} \rightarrow W^+W^-$, and $p\bar{p} \rightarrow Z^0 \rightarrow \tau^+\tau^-$ at a center-of-mass energy of $\sqrt{s} = 1.96$ TeV. We perform a likelihood fit to the dilepton data in a parameter space defined by the missing transverse energy and the number of jets in the event. Our results, which use 360 pb^{-1} of data recorded with the CDF II detector at the Fermilab Tevatron Collider, are $\sigma(t\bar{t}) = 8.5_{-2.2}^{+2.7} \text{ pb}$, $\sigma(W^+W^-) = 16.3_{-4.4}^{+5.2} \text{ pb}$, and $\sigma(Z^0 \rightarrow \tau^+\tau^-) = 291_{-46}^{+50} \text{ pb}$.

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There are relatively few standard model (SM) processes that contribute significantly to a final state containing a pair of highly energetic charged leptons. The processes that can contribute to these “dilepton” events include top-quark pair ($t\bar{t}$) production, W -boson pair (W^+W^-) production, and Drell-Yan processes. The distinctiveness of this final state offers the potential for unique tests of the SM and searches for new physics. In Run I (1992–1996) the Collider Detector at Fermilab (CDF) observed several $t\bar{t}$ candidate events in the dilepton decay mode [1] with unusual characteristics, and it was suggested that their kinematics could be better described by the cascade decays of heavy supersymmetric quarks [2]. Furthermore, the top-quark’s extraordinarily large mass might be an indication of a close connection with the mechanism of mass generation itself, as, for example, in the model of topcolor assisted technicolor [3], which predicts new resonances decaying to $t\bar{t}$. In addition, a fourth generation fermion family [4] could enhance the gluon-gluon Higgs production cross section by an order of magnitude, which, for a heavy Higgs boson, would lead to an increase in the number of W boson pairs observed [5]. These examples could all produce an excess of dilepton events in our data, which could affect the cross section measurements of the main SM processes. This provides the motivation behind the present study of highly energetic dilepton events.

The main SM processes with a high- p_T [6] dilepton final state can be identified based on their distinct event characteristics. For example, $t\bar{t}$ and W^+W^- production with decays in the $e\mu$ channel, $t\bar{t} \rightarrow W^+bW^-\bar{b} \rightarrow e^\pm\mu^\mp\nu\bar{\nu}b\bar{b}$ and $W^+W^- \rightarrow e^\pm\mu^\mp\nu\bar{\nu}$, and the di-tau decays of Z^0 bosons, $Z^0 \rightarrow \tau^+\tau^- \rightarrow e^\pm\mu^\mp\nu\nu\bar{\nu}\bar{\nu}$, can be distinguished from each other by the number of jets, N_j , and the missing transverse energy [7], $\cancel{E}_T$, in the event. Both $t\bar{t}$ and W^+W^- events typically have large $\cancel{E}_T$ from the final state undetected neutrinos; however, due to the two final-state b -quarks, $t\bar{t}$ has a greater number of jets than W^+W^- . Conversely, $Z^0 \rightarrow \tau^+\tau^-$ events have small $\cancel{E}_T$ (due to lower energy neutrinos typically traveling in opposite directions), and most often no jets. For both W^+W^- and $Z^0 \rightarrow \tau^+\tau^-$ events, jets arise only through higher-order processes that include initial-state gluon radiation.

In contrast to analyses dedicated to measuring a single SM process, the analysis presented here considers all events with a high- p_T electron and muon, without further selection requirements. We exploit the different $\cancel{E}_T$ and N_j characteristics to simultaneously extract the production cross sections of the three main processes described. This is done by fitting the $e\mu$ data in a two-dimensional (2-D) $\cancel{E}_T - N_j$ parameter space to template distributions of $t\bar{t}$, W^+W^- , and $Z^0 \rightarrow \tau^+\tau^-$ events. Less significant processes are also taken into account.

We also consider the ee and $\mu\mu$ final states, but in these cases we have the added complication of a large Drell-Yan ($Z^0/\gamma^* \rightarrow e^+e^-$ or $\mu^+\mu^-$) contribution, which would

overwhelm the $t\bar{t}$ and W^+W^- contributions without further requirements. We reduce the Drell-Yan contributions by requiring events to have significant $\cancel{E}_T$ in those channels. Our $t\bar{t}$ and W^+W^- results use all three dilepton final states. For extracting the $Z^0 \rightarrow \tau^+\tau^-$ cross section we use only $e\mu$ events as the removal of Drell-Yan ee and $\mu\mu$ events also significantly reduces the $Z^0 \rightarrow \tau^+\tau^-$ contribution in the ee and $\mu\mu$ channels.

Since this method makes minimal requirements on events after requiring two leptons, it utilizes the full statistical power of the data for given lepton definitions. In the present analysis we have chosen to use very tight lepton identification requirements to demonstrate the method, and establish a foundation for future measurements with more data. This gives us greater control of the background processes, in particular, those involving jets being misidentified as, or containing, leptons. However, it also reduces the potential statistical gain from the method. Even so, the results are comparable in precision to the analyses dedicated to the individual cross section measurements, and that use looser dilepton definitions. Furthermore, by looking at all processes simultaneously in the same generic dilepton sample, this method has the potential of testing the SM consistency differently from single cross section measurements. We also potentially have additional sensitivity to new physics which might fall into the $\cancel{E}_T - N_j$ parameter space that we study.

The data sample was collected with the CDF II detector between 2002 and 2004, and corresponds to an integrated luminosity of approximately 360 pb^{-1} . The results presented here complement the individual cross section measurements from CDF for each process [8–10], which make additional requirements to reduce, and assume the SM cross sections for all the processes in the dilepton sample other than the one being measured.

The CDF II detector has a general-purpose design [11]. It includes a tracking system inside a 1.4 T superconducting solenoidal magnet which is composed of silicon detectors for high-precision track measurements, and an open-cell drift chamber surrounding the silicon system. Electromagnetic and hadronic calorimeters surround the tracking system and solenoid and are used to measure the energy of interacting particles. In this analysis information is used from both the central (covering the pseudorapidity range $|\eta| < 1.1$) and end-plug detectors (of which we use the range $1.2 < |\eta| < 2.0$) to identify electron candidates. The missing transverse energy calculation uses calorimeter towers with $|\eta| < 3.5$. We identify jets using clusters of calorimeter towers above an energy threshold of 3 GeV, with fixed cone radius $\Delta R = \sqrt{\Delta\phi^2 + \Delta\eta^2} = 0.4$. The jet transverse energy is corrected for the calorimeter response and multiple interactions [12,13]. We require corrected jets to have $E_T > 15 \text{ GeV}$ and $|\eta| < 2.5$. A set of drift chambers located outside the central hadronic calorimeters are used to detect muons in the region $|\eta| < 0.6$. Additional

drift chambers and scintillation counters detect muons in the region $0.6 < |\eta| < 1.0$.

A three-level trigger system is used to select events. The triggers employed to collect events for this analysis are [14] an inclusive central electron ($|\eta| < 1.1$) trigger requiring an electron with $E_T > 18$ GeV, and an inclusive central muon ($|\eta| < 1.0$) trigger requiring a muon with $p_T > 18$ GeV/c.

Events selected contain two opposite sign leptons (electrons or muons) consistent with originating from the same vertex, and with $E_T > 20$ GeV for electrons and $p_T > 20$ GeV/c for muons. Muons from cosmic rays and electrons from photon conversions are removed [14]. Both leptons are required to be isolated in the calorimeter [15] and the tracking chamber [16], in order to reduce the probability that a jet is misidentified as a lepton, or that a selected lepton comes from the semileptonic decay of b or c hadrons.

We consider three classes of dilepton events: $e\mu$, ee , and $\mu\mu$. For the $e\mu$ channel we make no further requirements on the event after the selection of two leptons as described above, and simply count the number of jets, N_j , and measure the $\cancel{E}_T$ in the event.

The ee and $\mu\mu$ channels have a large Drell-Yan contribution which we reduce significantly by applying a requirement on the missing transverse energy significance defined by

$$\cancel{E}_T^{\text{sig}} = \frac{\cancel{E}_T}{\sqrt{\sum E_T}}$$

where $\sum E_T$ is the sum of transverse energies over all calorimeter towers, corrected to include the p_T of the muons. We require $\cancel{E}_T^{\text{sig}} > 2.5$ GeV $^{1/2}$ for all ee and $\mu\mu$ events. Without this cut our fit methodology is unable to separate the Drell-Yan contribution from the standard processes, in spite of different $\cancel{E}_T$ distribution shapes, due to the fact that the Drell-Yan background has more than 2 orders of magnitude more events than our signal processes.

After the above dilepton and $\cancel{E}_T^{\text{sig}}$ requirements, the dominant SM contributions which we consider as our signal processes, are $t\bar{t} \rightarrow W^+bW^-\bar{b} \rightarrow \ell^+\nu b\ell^-\bar{\nu}\bar{b}$, $W^+W^- \rightarrow \ell^+\nu\ell^-\bar{\nu}$, and $Z^0 \rightarrow \tau^+\tau^- \rightarrow \ell^+\nu_\ell\bar{\nu}_\ell\ell^-\bar{\nu}_\ell\nu_\tau$, where ℓ indicates an electron or muon. The predicted distributions of these signal processes in the $\cancel{E}_T - N_j$ parameter space for the $e\mu$ channel are shown in Fig. 1, and were obtained from Monte Carlo simulations.

Background processes, include Drell-Yan (ee , $\mu\mu$), WZ , ZZ , $W + \gamma$, and $W + \text{jets}$. The predicted combined background distributions in the $\cancel{E}_T - N_j$ parameter space is also shown in Fig. 1. We fit the data $e\mu$, ee , and $\mu\mu$ distributions to the expected signal shapes, letting each of their normalizations float in the fit. The $e\mu$ data distribution is shown in Fig. 2, of which about 60% of the events are expected to be from $Z^0 \rightarrow \tau^+\tau^-$ and concentrated in

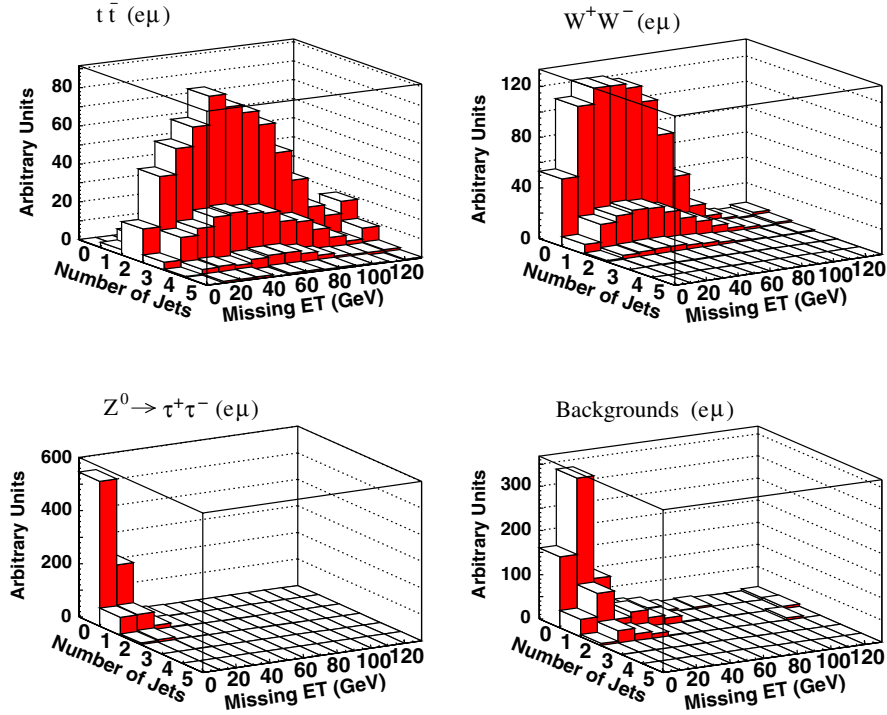


FIG. 1 (color online). The $\cancel{E}_T - N_j$ distributions for $t\bar{t}$, W^+W^- , $Z^0 \rightarrow \tau^+\tau^-$ in the $e\mu$ channel. Also shown is the combined background distribution for the $e\mu$ channel. All the distributions are normalized to an arbitrary equal volume. The numbers of events expected in 360 pb^{-1} for each source are given in Table III. The highest $\cancel{E}_T$ and N_j bins include any overflow events.

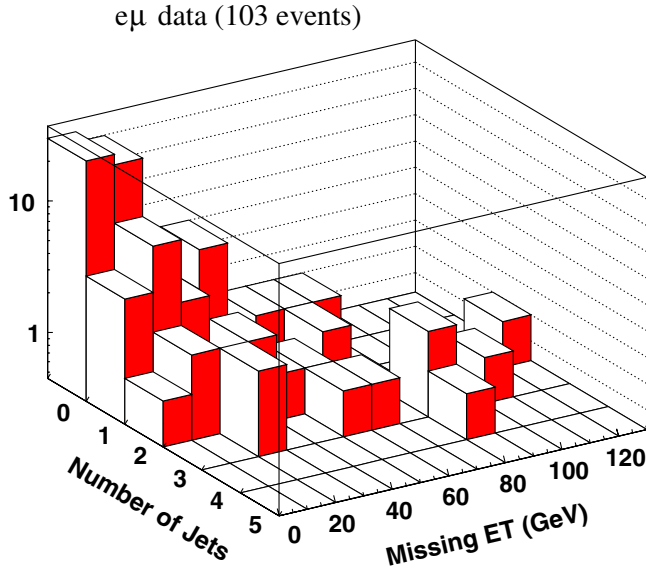


FIG. 2 (color online). The $\cancel{E}_T - N_j$ distribution for data in the $e\mu$ channel. The log scale is used to make the low count bins more visible.

the low- $\cancel{E}_T$, zero-jet region of the $\cancel{E}_T - N_j$ parameter space.

The acceptances for our signal processes are shown in Table I, calculated using the PYTHIA [17] Monte Carlo program, followed by a full simulation of the CDF II detector which is based on the GEANT simulation program [18]. The $Z^0 \rightarrow \tau^+ \tau^-$ acceptances are defined as the fraction of $Z^0/\gamma^* \rightarrow \tau^+ \tau^-$ events generated in the di-tau mass range of $66 < M_{\tau\tau} < 116 \text{ GeV}/c^2$ that pass our dilepton selection criteria. The contribution from γ^* is only about 0.3%. For the ee and $\mu\mu$ channels the additional $\cancel{E}_T^{\text{sig}}$ requirement removes 32% of $t\bar{t}$ events and 36% of

W^+W^- events, while also removing about 99.8% of Drell-Yan events.

The systematic uncertainties on the $t\bar{t}$, W^+W^- , and $Z^0 \rightarrow \tau^+ \tau^-$ acceptances, summarized in Table II, result from uncertainties in the jet energy scale (JES) [13], the modeling of the initial-state radiation (ISR) by PYTHIA [and for the case of $t\bar{t}$ also final-state radiation (FSR)] [19], the uncertainty on the parton distribution functions [20], the modeling of $\cancel{E}_T^{\text{sig}}$, and uncertainties in the lepton trigger and identification efficiencies. The $\cancel{E}_T^{\text{sig}}$ systematic uncertainty does not apply to the $Z^0 \rightarrow \tau^+ \tau^-$ cross section measurement, as we fit for this only in the $e\mu$ channel. Moreover, the $e\mu$ channel acceptance does not suffer from a jet energy scale uncertainty, as we include $e\mu$ events with all jet multiplicities, whereas for the ee and $\mu\mu$ channels this systematic uncertainty enters through the $\cancel{E}_T^{\text{sig}}$ requirement. A 6% uncertainty on the integrated luminosity is applied to the expected number of events for all processes [21].

The Drell-Yan and diboson backgrounds are determined using PYTHIA Monte Carlo, followed by the detector simulation. We normalize the total number of events for these processes to theoretical cross section predictions [22]. To estimate the $W + \gamma$ background we use a matrix element generator [23] and use PYTHIA for the initial-state QCD radiation and hadronization. The $W + \text{jets}$ background, where a jet or track is misidentified as an electron or muon, is determined from the data. We calculate the probability, using inclusive jet samples and corrected for sources of real leptons, that a jet with a large fraction of its energy deposited in the electromagnetic calorimeter is misidentified as an electron, and the probability that a minimum ionizing track is misidentified as a muon. The background is then determined by applying these probabilities to jets in a data sample of $(W \rightarrow \ell\nu) + \text{jets}$ events.

TABLE I. Summary of acceptances for $t\bar{t}$, W^+W^- , and $Z^0 \rightarrow \tau^+ \tau^-$ events, where the quoted errors include the systematic uncertainties from Table II. Values include SM branching fractions to the dilepton final state. The $t\bar{t}$ events were simulated with a top-quark mass of $178 \text{ GeV}/c^2$.

	$e\mu$	ee	$\mu\mu$
$t\bar{t}$	$(0.399 \pm 0.029)\%$	$(0.144 \pm 0.019)\%$	$(0.136 \pm 0.015)\%$
W^+W^-	$(0.294 \pm 0.018)\%$	$(0.111 \pm 0.008)\%$	$(0.092 \pm 0.006)\%$
$Z \rightarrow \tau\tau$	$(0.0458 \pm 0.0032)\%$	$(0.0008 \pm 0.0001)\%$	$(0.0005 \pm 0.0001)\%$

TABLE II. Summary of systematic uncertainties on the acceptance for each “signal” process. See text for further details.

Source	$t\bar{t}(ee)$	$t\bar{t}(e\mu)$	$t\bar{t}(\mu\mu)$	$W^+W^-(ee)$	$W^+W^-(e\mu)$	$W^+W^-(\mu\mu)$	$Z^0 \rightarrow \tau^+ \tau^-(e\mu)$
JES	5%	—	6%	1%	—	1%	—
ISR	8%	4%	6%	5%	5%	5%	5%
FSR	7%	3%	5%	—	—	—	—
Other	6%	5%	5%	5%	4%	4%	5%
Total	13%	7%	11%	7%	6%	7%	7%

As a result of very low statistics and other large uncertainties in this calculation, we assume a 100% total uncertainty on our final W + jets background estimates. A summary of all expected contributions for each dilepton channel is given in Table III, together with observed numbers of events.

We extract the $t\bar{t}$, W^+W^- , and $Z^0 \rightarrow \tau^+\tau^-$ cross sections simultaneously for the $e\mu$ channel by maximizing the binned likelihood function:

$$L(\mu_{t\bar{t}}, \mu_{WW}, \mu_{\tau\tau}) = \prod_i^{N_{\text{bin}}} \frac{\mu_i^{n_i} e^{-\mu_i}}{n_i!} \times \prod_j G(x_j, \sigma_j) \quad (1)$$

where the index i runs over all N_{bin} bins in the two-dimensional $\cancel{E}_T - N_j$ parameter space, and j runs over all variables x_j , which are parameters in the likelihood function constrained by a Gaussian of width σ_j , the estimated uncertainty on x_j . These variables consist of the expected number of events for the background processes (given in Table III), the acceptances for the signal processes (given in Table I), and the integrated luminosity.

The parameters $\mu_{t\bar{t}}$, μ_{WW} , and $\mu_{\tau\tau}$ are the expected total numbers of $t\bar{t}$, W^+W^- , and $Z^0 \rightarrow \tau^+\tau^-$ events, respectively. The expected distributions of these processes in the $\cancel{E}_T - N_j$ parameter space determine the probabilities $p_{t\bar{t},i}$, $p_{WW,i}$, and $p_{\tau\tau,i}$ that a $t\bar{t}$, W^+W^- , or $Z^0 \rightarrow \tau^+\tau^-$ event, respectively, will appear in the i th bin. Therefore, the total expected number of events in the i th bin of the $\cancel{E}_T - N_j$ parameter space is:

$$\mu_i = (p_{t\bar{t},i} \times \mu_{t\bar{t}}) + (p_{WW,i} \times \mu_{WW}) + (p_{\tau\tau,i} \times \mu_{\tau\tau}) + n_{\text{other}} \quad (2)$$

where n_{other} is the expected number of events in the i th bin

from all background processes, and is fixed in the likelihood fit within its Gaussian constraint. The product over i in the likelihood function is the product of Poisson probabilities for each bin, in which the expected number of events is μ_i , and n_i is the corresponding number of events observed in the data. We use 10 GeV wide bins for $\cancel{E}_T < 60$ GeV, and 20 GeV bins for $60 < \cancel{E}_T < 120$ GeV. Note that in Figs. 1 and 2 we show the $\cancel{E}_T - N_j$ shapes using a uniform 10 GeV binning in $\cancel{E}_T$.

For each of the signal processes, $\mu_k = \sigma_k \epsilon_k \mathcal{L}$, where $k = t\bar{t}$, W^+W^- , or $Z^0 \rightarrow \tau^+\tau^-$, and σ_k is the production cross section for the process k , which is a free parameter in the likelihood fit. The acceptances, ϵ_k , and integrated luminosity of the data sample, $\mathcal{L}$, are fixed within a Gaussian constraint as mentioned above. By maximizing the likelihood function we extract the production cross sections for $t\bar{t}$, W^+W^- , and $Z^0 \rightarrow \tau^+\tau^-$.

We perform a similar likelihood fit to the full $ee + e\mu + \mu\mu$ data to extract $t\bar{t}$ and W^+W^- cross sections, in which we consider $Z^0 \rightarrow \tau^+\tau^-$ as a fixed background included in the n_{other} term, since in the ee and $\mu\mu$ channels it has been significantly reduced by the $\cancel{E}_T^{\text{sig}}$ requirement. We perform the full fit using a likelihood function that is the product of the individual likelihood functions for each channel. For the full fit we also assume that within a given channel (ee , $e\mu$, or $\mu\mu$) the signal acceptances are 100% correlated since they are driven by the lepton identification efficiencies, and that there are no correlations between different channels. The latter is not quite true because, for example, the ee and $e\mu$ channels have correlations in the lepton identification uncertainties due to overlap in lepton types. However, by varying these correlations between 0% and 100% we see a negligible effect on the extracted cross sections.

TABLE III. The numbers of SM predicted and observed events, in 360 pb^{-1} of data. For the $e\mu$ channel, where the only requirement is two high- p_T leptons, the first three processes are considered signals for which cross sections are measured. For the ee and $\mu\mu$ channels, where an additional $\cancel{E}_T^{\text{sig}}$ requirement is made, only the first two processes are regarded as signals. To calculate the expected number of events from our signal processes we used the cross section central values, $\sigma(t\bar{t}) = 6.1 \text{ pb}$, $\sigma(W^+W^-) = 12.4 \text{ pb}$, and $\sigma(Z^0 \rightarrow \tau^+\tau^-) = 251 \text{ pb}$. Uncertainties on the theoretical cross sections are not included.

	$e\mu$	ee	$\mu\mu$	$\ell\ell$
$t\bar{t}$	10.0 ± 0.7	3.6 ± 0.5	3.4 ± 0.4	17.0 ± 1.6
W^+W^-	13.8 ± 0.8	5.2 ± 0.4	4.3 ± 0.3	23.3 ± 1.5
$Z^0 \rightarrow \tau^+\tau^-$	57.8 ± 4	1.1 ± 0.2	0.6 ± 0.1	59.5 ± 4.3
$DY \rightarrow ee$	0	15.4 ± 3.2	0	15.4 ± 3.2
$DY \rightarrow \mu\mu$	9.3 ± 0.8	0	11.6 ± 2.4	20.8 ± 3.2
WZ	0.70 ± 0.06	1.26 ± 0.09	1.11 ± 0.08	3.07 ± 0.23
ZZ	0.07 ± 0.01	0.47 ± 0.03	0.42 ± 0.03	0.96 ± 0.07
$W\gamma$	1.2 ± 0.5	1.8 ± 0.7	0	3.0 ± 1.2
W + jets	3.0 ± 3.0	2.1 ± 2.1	1.6 ± 1.6	6.8 ± 6.8
Total SM	96 ± 5	31 ± 4	23 ± 3	150 ± 12
Data	103	24	29	156

TABLE IV. Cross section measurements from a global fit of 360 pb^{-1} of high- p_T dilepton data. The first uncertainty is the error returned by the likelihood fit, which includes statistical, acceptance systematic, and luminosity uncertainties. The magnitudes of the latter two uncertainties are given in Table II and in the text. The second is the systematic uncertainty in the template shapes.

Process	$e\mu$	$ee + \mu\mu + e\mu$
$\sigma(t\bar{t})$	$9.3^{+3.1+0.7}_{-2.6-0.2} \text{ pb}$	$8.5^{+2.6+0.7}_{-2.2-0.3} \text{ pb}$
$\sigma(W^+W^-)$	$11.4^{+5.2+0.5}_{-4.3-0.1} \text{ pb}$	$16.3^{+5.1+0.8}_{-4.4-0.2} \text{ pb}$
$\sigma(Z^0 \rightarrow \tau^+\tau^-)$	$291^{+50+6}_{-46-3} \text{ pb}$	—

For both the $e\mu$ -only and $ee + e\mu + \mu\mu$ scenarios we perform two sets of fits. Our main results are obtained by letting all the signal cross sections float, with the exception of $Z^0 \rightarrow \tau^+\tau^-$ in the ee and $\mu\mu$ channels, to simultaneously extract the signal cross sections from the fit. These results are summarized in Table IV. As a result of the good separation of our signal processes in the $\cancel{E}_T - N_j$ parameter space, we observe very little correlation between these cross section measurements. In the $e\mu$ channel fit where all three cross sections float, these correlations are about -0.06 between $\sigma(t\bar{t})$ and $\sigma(W^+W^-)$, -0.05 between $\sigma(t\bar{t})$ and $\sigma(Z^0 \rightarrow \tau^+\tau^-)$, and, -0.19 between $\sigma(W^+W^-)$ and $\sigma(Z^0 \rightarrow \tau^+\tau^-)$. The $t\bar{t}$ production cross section measurement is relatively insensitive to the top mass used in generating the $\cancel{E}_T - N_j$ template shape. We generated templates using a top mass of $178 \text{ GeV}/c^2$ [24]. Using simulated experiments we observe a 1% variation in $\sigma(t\bar{t})$ if the top mass is varied between $165 \text{ GeV}/c^2$ and $178 \text{ GeV}/c^2$; therefore, we neglect any effects due to the uncertainty on the top mass.

Fits to the data are also performed by fixing all but one of the signal processes to their SM expected values (we use $\sigma(t\bar{t}) = 6.1 \pm 0.9 \text{ pb}$ [25], $\sigma(W^+W^-) = 12.4 \pm 0.8 \text{ pb}$ [26], and $\sigma(Z^0 \rightarrow \tau^+\tau^-) = 251.3 \pm 5.0 \text{ pb}$ [27]). We con-

sider the results extracted from these fits as cross-checks, as they use the added constraint of assuming SM production for all processes other than the one being measured, in a similar fashion to the more standard counting experiment results. Significant differences between these cross-checks and our main results could be an indication that the SM assumptions being used are incorrect. We summarize the results from these fits in Table V. We note that the W^+W^- cross section from $e\mu$ events differs from the full fit result indicating an excess in the $ee + \mu\mu$ channels. This is consistent with the observation from the dedicated W^+W^- measurement [9].

The uncertainties on the measured cross sections include a component coming from the fit (including statistical, acceptance systematic, and integrated luminosity) and a second one due to changes in the $\cancel{E}_T^{\text{sig}} - N_j$ distributions caused by the systematic sources mentioned in Table II. The fit program used was MINUIT [28], which minimizes $-\ln(L)$, and determines the fit errors from the $-2\ln(L)$ values. To evaluate systematic changes in the shape of $\cancel{E}_T - N_j$ distributions, we use simulated experiments with one of the signal or background distributions from a Monte Carlo simulation with a particular systematic effect applied. The shape systematic uncertainties are summarized in Table VI.

In summary, we present a new method to study globally the production of events with final states that include two high- p_T leptons. We measure simultaneously the production cross sections for $t\bar{t}$, W^+W^- , and $Z^0 \rightarrow \tau^+\tau^-$ and find they are in good agreement with SM theoretical predictions. In addition to the potential this analysis technique has for precision cross section measurements in the dilepton channel with more data, it could also be promising for model independent searches for new physics.

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TABLE V. Cross section measurements from a fit to the data with all but one signal process fixed to its SM value. The uncertainties have the same meaning as in Table IV.

Process	$e\mu$	$ee + \mu\mu + e\mu$
$\sigma(t\bar{t})$ (W^+W^- , $Z^0 \rightarrow \tau^+\tau^-$ fixed)	$9.3^{+3.1+0.7}_{-2.6-0.2} \text{ pb}$	$8.4^{+2.5+0.7}_{-2.1-0.3} \text{ pb}$
$\sigma(W^+W^-)$ ($t\bar{t}$, $Z^0 \rightarrow \tau^+\tau^-$ fixed)	$12.3^{+5.3+0.5}_{-4.4-0.1} \text{ pb}$	$16.1^{+5.0+0.8}_{-4.3-0.2} \text{ pb}$
$\sigma(Z \rightarrow \tau\tau)$ ($t\bar{t}$, W^+W^- fixed)	$293^{+49+6}_{-45-3} \text{ pb}$	—

TABLE VI. Summary of the shape systematic uncertainties for the $e\mu$ and full ($ee + e\mu + \mu\mu$) fit.

Source	$t\bar{t}(e\mu)$	$W^+W^-(e\mu)$	$Z^0 \rightarrow \tau^+\tau^-(e\mu)$	$t\bar{t}$ (full)	W^+W^- (full)
JES	$+6\%$ -1%	$+4\%$ -1%	$+2\%$ -1%	$+7\%$ -2%	$+5\%$ -1%
ISR	$+4\%$ -2%	$\pm 1\%$	$\pm 1\%$	$+5\%$ -2%	$\pm 1\%$
FSR	$\pm 1\%$	—	—	$\pm 1\%$	—
Total	$+7\%$ -2%	$+4\%$ -1%	$+2\%$ -1%	$+8\%$ -3%	$+5\%$ -1%

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