

Comments and Addenda

The Comments and Addenda section is for short communications which are not of such urgency as to justify publication in Physical Review Letters and are not appropriate for regular articles. It includes only the following types of communications: (1) comments on papers previously published in The Physical Review or Physical Review Letters; (2) addenda to papers previously published in The Physical Review or Physical Review Letters, in which the additional information can be presented without the need for writing a complete article. Manuscripts intended for this section may be accompanied by a brief abstract for information retrieval purposes. Accepted manuscripts will follow the same publication schedule as articles in this journal, and galley proofs will be sent to authors.

Comment on "Threshold Kinematic Constraints and Correlations of Helicity Amplitudes"

JERROLD FRANKLIN

Department of Physics, Temple University, Philadelphia, Pennsylvania 19122

(Received 15 April 1970)

Previously derived threshold constraints on helicity amplitudes are simplified for high-spin cases, using a suggestion made by King and Kuo.

IN a recent paper,¹ we used partial-wave threshold behavior² to derive threshold (and pseudothreshold) constraints on helicity amplitudes. These relations appeared to require derivative (with respect to t) constraints for some cases when the total spin (S) was ≥ 2 . On the other hand, the results, for similar cases, of Cohen-Tannoudji *et al.*³ and of King and Kuo⁴ do not involve derivative constraints. The purpose of this comment is to reconcile this apparent difference by noting that the presence of the derivatives in our original derivation was unnecessary. The derivatives appeared to be required because the helicity (λ) dependence could not be factored out of the partial-wave sum (on J) for some cases with $S \geq 2$. This was because the J, λ dependence in Eq. (40) of I appeared in a nonfactorizable polynomial $P_N(S, \lambda; J)$. However, it turns out that only the leading power of λ in $P_N(S, \lambda; J)$ need be considered for the constraint equation. This follows immediately from the simple observation of King and Kuo⁴ that lower-order terms in λ would be removed by the lower-order constraints. The remaining J, λ dependence is obviously factori-

zable. The result is that derivatives are unnecessary in any of the constraint relations, and the complicated operator $P_N(S, \lambda; J_{op})$ can be replaced by simply λ^N in the general constraint equations (45)–(48) of I. These equations can then be written¹

$$\lim_{W \rightarrow m+\mu} q^{-S+N} \sum_{\lambda} (-1)^{-\lambda/2} [(S+\lambda)! (S-\lambda)!]^{-1/2} \\ \times \lambda^N A_{S'\lambda'; S\lambda}(W, t) = \alpha_{S'\lambda'; S^{(N)}}(t), \quad N \leq S \quad (45')$$

$$\lim_{W \rightarrow m'+\mu'} q'^{-S'+N'} \sum_{\lambda'} (-1)^{\lambda'/2} [(S'+\lambda')! (S'-\lambda')!]^{-1/2} \\ \times \lambda'^{N'} A_{S'\lambda'; S\lambda}(W, t) = \alpha'_{S'\lambda'; S\lambda^{(N')}}(t), \quad N' \leq S' \quad (47')$$

for inelastic scattering, and

$$\lim_{W \rightarrow m+\mu} q^{-S-S'+N+N'} \\ \times \sum_{\lambda, \lambda'} \frac{(-1)^{(\lambda'-\lambda)/2} \lambda^N \lambda'^{N'}}{[(S+\lambda)! (S-\lambda)! (S'+\lambda')! (S'-\lambda')!]^{1/2}} \\ \times A_{S'\lambda'; S\lambda}(W, t) = \alpha_{S'\lambda'; S^{(N, N')}}(t), \quad N \leq S, \quad N' \leq S' \\ (-1)^{S-S'-N+N'} = \eta \text{ (relative intrinsic parity)} \quad (48')$$

for elastic scattering. These agree with the equivalent results of King and Kuo.⁴

¹ J. Franklin, Phys. Rev. **170**, 1606 (1968).

² J. Franklin, Phys. Rev. **152**, 1437 (1966).

³ G. Cohen-Tannoudji, A. Morel, and H. Navelet, Ann. Phys. (N. Y.) **46**, 239 (1968).

⁴ M. J. King and P. K. Kuo, Phys. Rev. D **1**, 442 (1970).