

Comments and Addenda

The Comments and Addenda section is for short communications which are not of such urgency as to justify publication in Physical Review Letters and are not appropriate for regular articles. It includes only the following types of communications: (1) comments on papers previously published in The Physical Review or Physical Review Letters; (2) addenda to papers previously published in The Physical Review or Physical Review Letters, in which the additional information can be presented without the need for writing a complete article. Manuscripts intended for this section may be accompanied by a brief abstract for information retrieval purposes. Accepted manuscripts will follow the same publication schedule as articles in this journal, and galley proofs will be sent to authors.

Experimental Limit on $\nu_\mu + e^- \rightarrow \nu_\mu + e^-$ and the Existence of Weakly Coupled Neutral Lepton Currents

CARL H. ALBRIGHT

Department of Physics, Northern Illinois University, DeKalb, Illinois 60115

(Received 20 April 1970)

A recent analysis by Steiner of the CERN neutrino spark-chamber experiment is employed to derive an upper limit on the elastic cross section for $\nu_\mu - e^-$ scattering. The limit obtained is just at a level consistent with the theoretical cross section predicted for neutral lepton currents coupled to form a weak isospin-invariant Hamiltonian.

RECENTLY, Steiner¹ has placed an upper limit of

$$\sigma_{\text{expt}} < 40\sigma_{V-A} \quad (1)$$

on the cross section for

$$\nu_e + e^- \rightarrow \nu_e + e^-, \quad (2)$$

from an analysis of the 1963-64 CERN neutrino spark-chamber experiment.² This result is of great interest, since a recently reported result³ indicates an anomalously large cross section for the related scattering of electrons by electron-type antineutrinos. The purpose of this note is to point out that an upper limit can also be placed on the cross section for

$$\nu_\mu + e^- \rightarrow \nu_\mu + e^- \quad (3)$$

by the same type of analysis. The limit here is of interest, for (3) can proceed in lowest order directly only through weakly coupled neutral currents by virtue of muon-number conservation or through weakly coupled charged currents mediated by second-order electromagnetic interactions. The result derived below for (3) is, in fact, much more restrictive than that given in (1) since the muon-type neutrino flux is considerably

larger than the electron-type neutrino flux as shown⁴ in Fig. 1.

The analysis of Steiner¹ indicates that of the 30 electron events observed with $E_e > 1$ GeV, which are compatible with the quasielastic process

$$\nu_e + n \rightarrow p + e^-, \quad (4)$$

one event may possibly be identified with process (2). From the universal $V-A$ charged-current theory⁵ and a knowledge of the detection efficiency of the experimental setup and the ν_e flux from the K_{e3} decays, Steiner finds that 0.11 ± 0.04 events of type (2) were expected. Poisson statistics then lead him to the upper limit on the experimental cross section quoted in (1) with 90% confidence.

We now note that the one event in 30 singled out is also compatible with process (3), so the same argument is applicable here. Comparison of the neutrino fluxes in Fig. 1 leads to an average ν_μ flux 120 times greater than the average ν_e flux for $E_\nu > 1$ GeV and to a rate for (3) 100 times that for (2) if the matrix elements for the two processes are assumed to have the same form.⁶ Hence 11 ± 4 events would have been expected, whereas only at most one was seen. With 90% confidence, we

¹ H. J. Steiner, Phys. Rev. Letters **24**, 746 (1970).

² J. K. Bienlein, A. Bohm, G. von Dardel, H. Faissner, F. Ferrero, J.-M. Gaillard, H. J. Gerber, B. Hahn, V. Kaftanov, F. Krienen, M. Reinharz, R. A. Salmeron, P. G. Seiler, A. Staude, J. Stein, and H. J. Steiner, Phys. Letters **13**, 80 (1964). See also Ref. 1, footnotes 4 and 6 therein.

³ F. Reines and H. S. Gurr, University of California at Irvine Report No. UCI 10, pp. 19-28 (unpublished). A revised version of this report has been published by F. Reines and H. Gurr, Phys. Rev. Letters **24**, 1448 (1970), in which a new upper limit of $\sigma_{\text{expt}} < 40\sigma_{V-A}$ is quoted for the antineutrino process corresponding to (2).

⁴ The neutrino fluxes pictured in Fig. 1 are obtained from Ref. 1 and M. Holder, A. Staude, A. Bohm, H. Faissner, H. J. Steiner, J. K. Bienlein, G. von Dardel, F. Ferrero, J.-M. Gaillard, H. J. Gerber, V. Kaftanov, F. Krienen, C. Manfredotti, M. Reinharz, and R. A. Salmeron, Nuovo Cimento **57A**, 338 (1968).

⁵ R. P. Feynman and M. Gell-Mann, Phys. Rev. **109**, 193 (1958); E. C. G. Sudarshan and R. E. Marshak, *ibid.* **109**, 1860 (1958).

⁶ On the basis of a contact interaction between the two lepton currents, one expects a cross section which rises linearly with the incident neutrino energy.

can then place the upper limit

$$\sigma_{\text{expt}}(\nu_\mu + e^- \rightarrow \nu_\mu + e^-) < 0.40\sigma_{V-A}. \quad (5)$$

It is interesting to speculate what kinds of neutral lepton currents may exist. For example, if the neutral currents couple only charged leptons (electrons or muons) or only neutrinos, elastic $\nu_\mu - e^-$ scattering will occur in lowest-order weak interactions only through terms of order G^2 . On the other hand, it is reasonable to guess that the charged and neutral lepton currents couple so as to form an isospin-invariant Hamiltonian according to

$$H_W^I = (G/\sqrt{2}) \sum_{l, l' = e, \mu} [l_\lambda^{(+)} l_\lambda'^{(-)} + l_\lambda^{(3)} l_\lambda'^{(3)}], \quad (6)$$

where the lepton current is just $l_\lambda^{(k)} = \frac{1}{2} \bar{\psi}_l \tau^{(k)} \gamma_\lambda (1 + \gamma_5) \psi_l$, and ψ_l is a two-component isospinor. The first term in (6) contains the charged-current contribution to muon decay, while the second term contributes to the process (3). On the basis of this model, $\sigma_{\text{th}}(\nu_\mu + e^- \rightarrow \nu_\mu + e^-) = \frac{1}{4}\sigma_{V-A}$, and the limit in (5) becomes

$$\sigma_{\text{expt}}(\nu_\mu + e^- \rightarrow \nu_\mu + e^-) < 1.6\sigma_{\text{th}} \quad (5')$$

with 90% confidence. This model is quite consistent⁷ with the present information.

If neutral currents exist and (6) is correct, one should note that elastic $\nu_e - e^-$ scattering can occur by both charged and neutral currents. The weak leptonic Hamiltonian of (6) predicts that $\sigma_{\text{th}}(\nu_e + e^- \rightarrow \nu_e + e^-) = \frac{1}{4}\sigma_{V-A}$. Hence if we reattribute the one possible electron event to (2), the upper limit (1) can be restated as

$$\sigma_{\text{expt}}(\nu_e + e^- \rightarrow \nu_e + e^-) < 160\sigma_{\text{th}} \quad (1')$$

with 90% confidence. On the other hand, if the one event in question should be identified with (3) rather

⁷ We should point out that if it were known that no events of process (3) have really been observed, the upper limit in (5') would become $\sigma_{\text{expt}} < 1.0\sigma_{\text{th}}$ with 90% confidence.

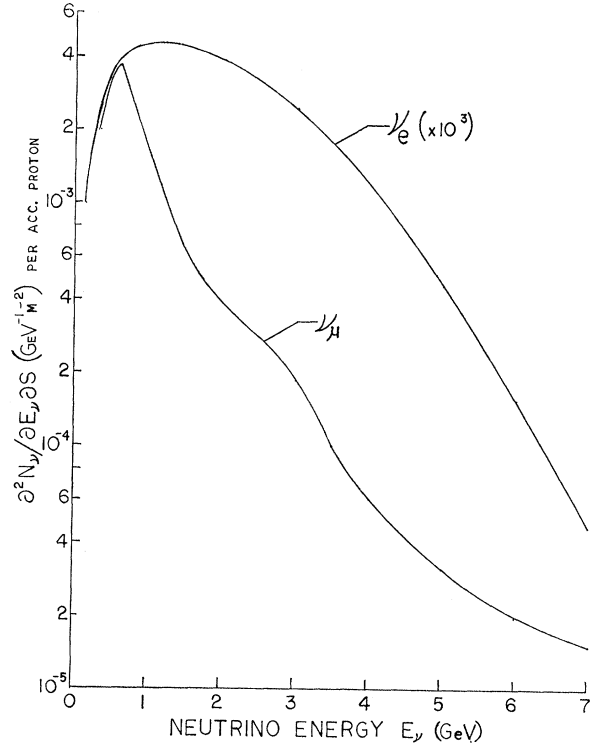


FIG. 1. Energy spectra of muon-type neutrinos and electron-type neutrinos in the CERN neutrino spark-chamber experiment.

than (2), the above is reduced to

$$\sigma_{\text{expt}}(\nu_e + e^- \rightarrow \nu_e + e^-) < 100\sigma_{\text{th}}. \quad (1'')$$

It is rather remarkable that this analysis indicates that the present experimental upper limits are within one order of magnitude of testing the predictions of the weak isospin-invariant Hamiltonian stated in (6). The next generation of neutrino experiments should clarify the matter decisively in the near future.

The author gratefully acknowledges Dr. M. Derrick for the very helpful conversations concerning this analysis held with him at Argonne National Laboratory. He also wishes to thank Dr. L. Desrosiers and Dr. P. J. O'Donnell for pointing out an error in the original manuscript.